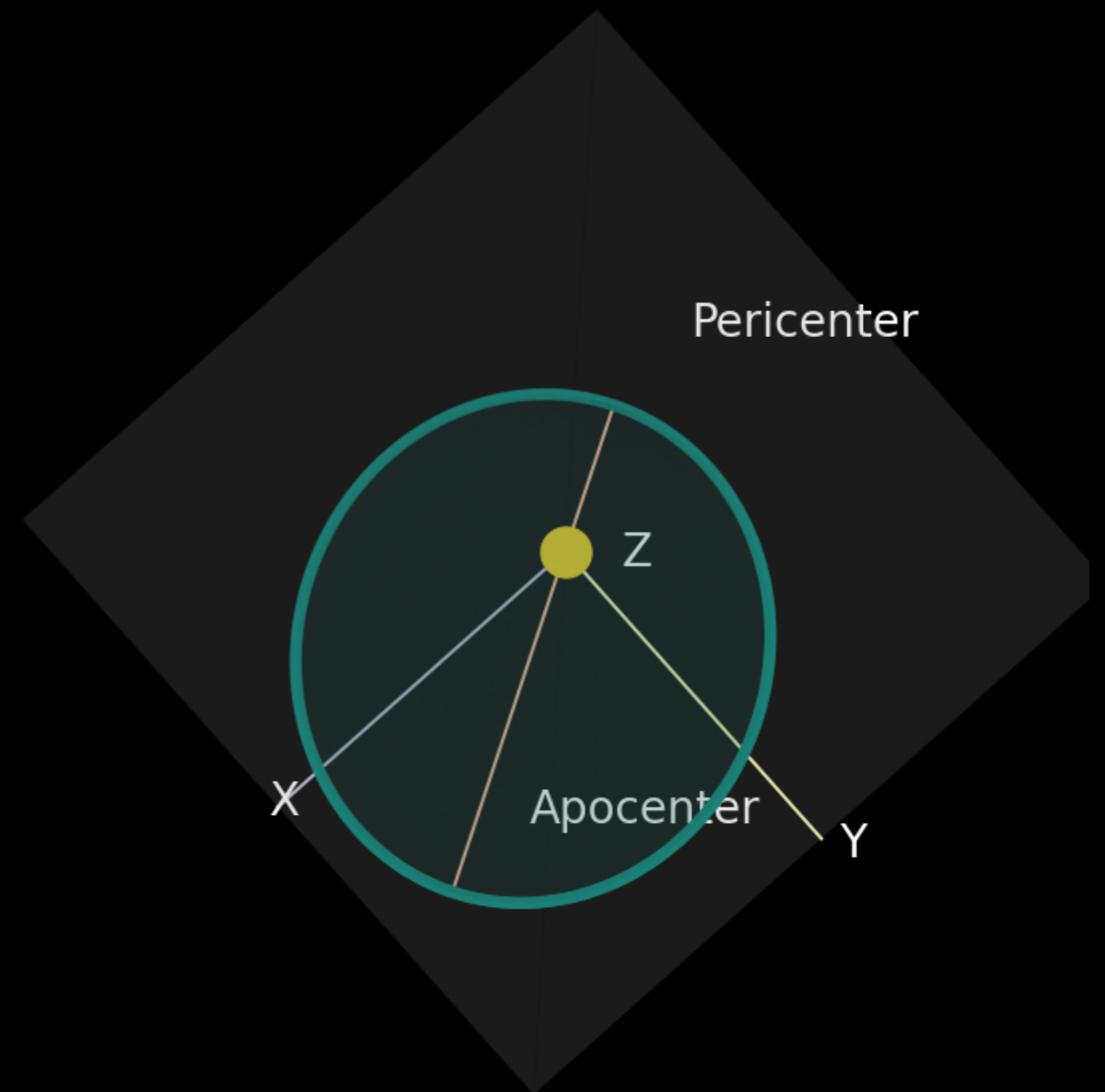


Evolution of Stellar structures in the Galactic Centre

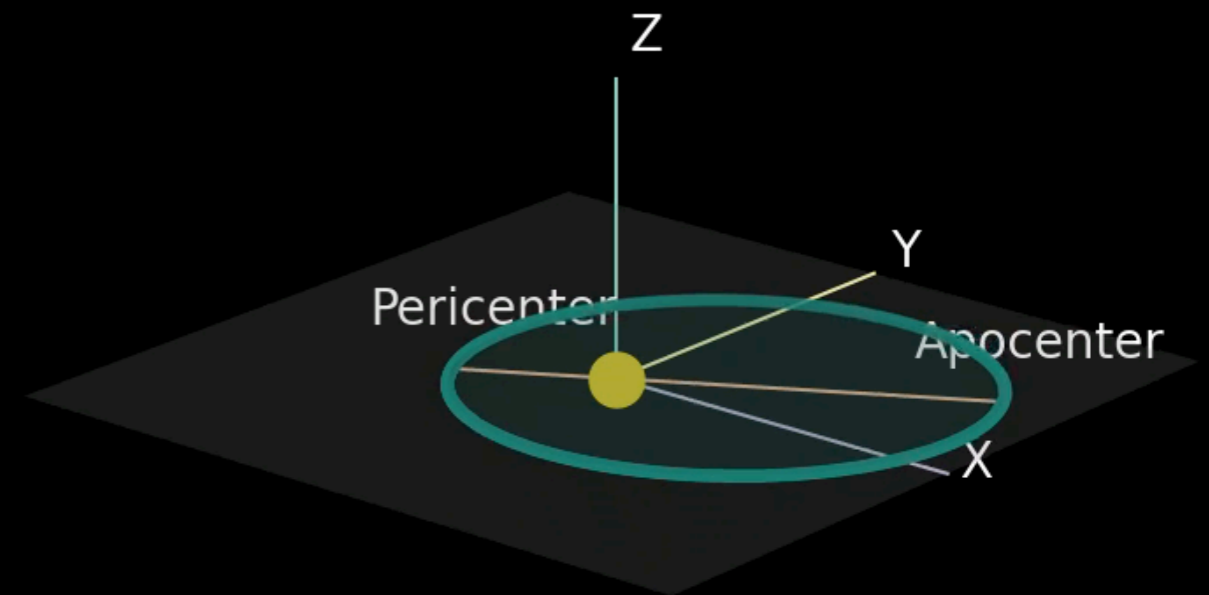
Keplerian Orbital Elements

- Shape of orbit:
 - Semi-major axis (a)
 - Eccentricity (e)
- Orientation of plane:
 - Inclination (i)
 - Longitude of ascending node (Ω)
- Orientation of particle
 - Argument of pericenter (ω)
 - True anomaly (ν)



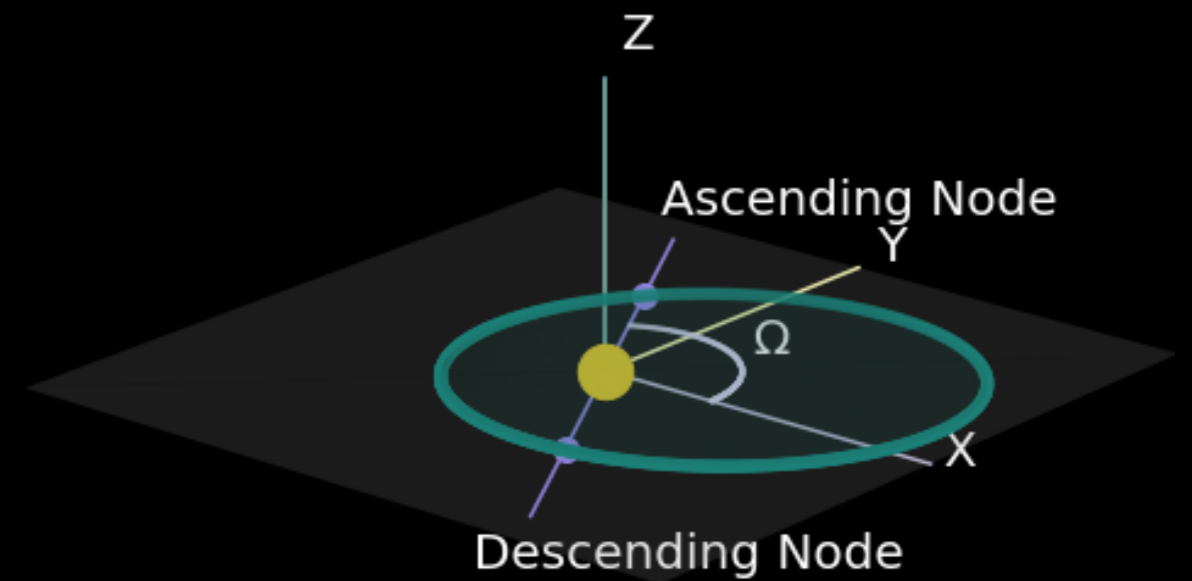
Keplerian Orbital Elements

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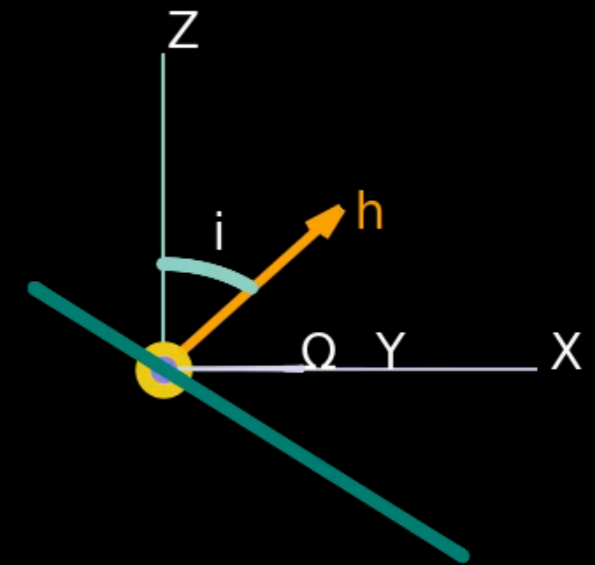
Keplerian Orbital Elements

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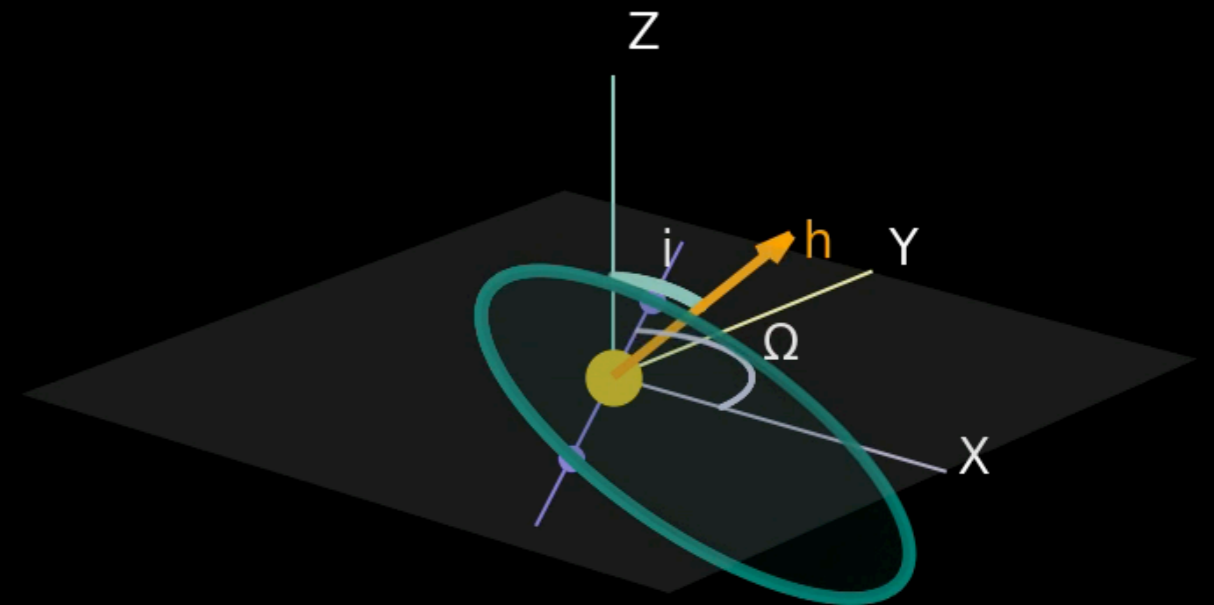
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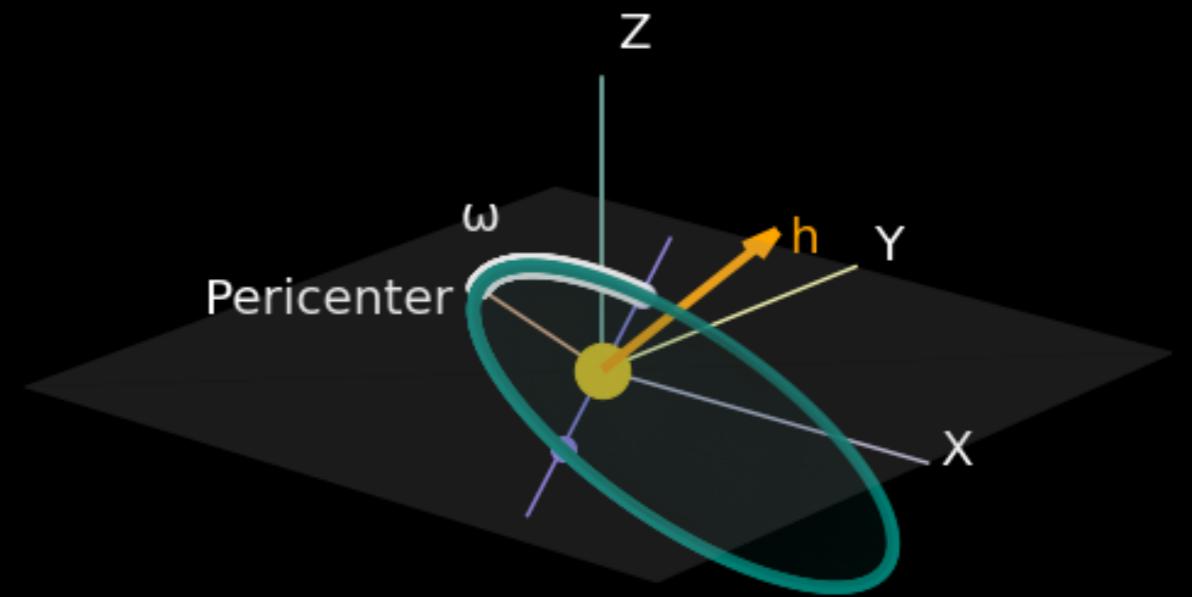
Keplerian Orbital Elements

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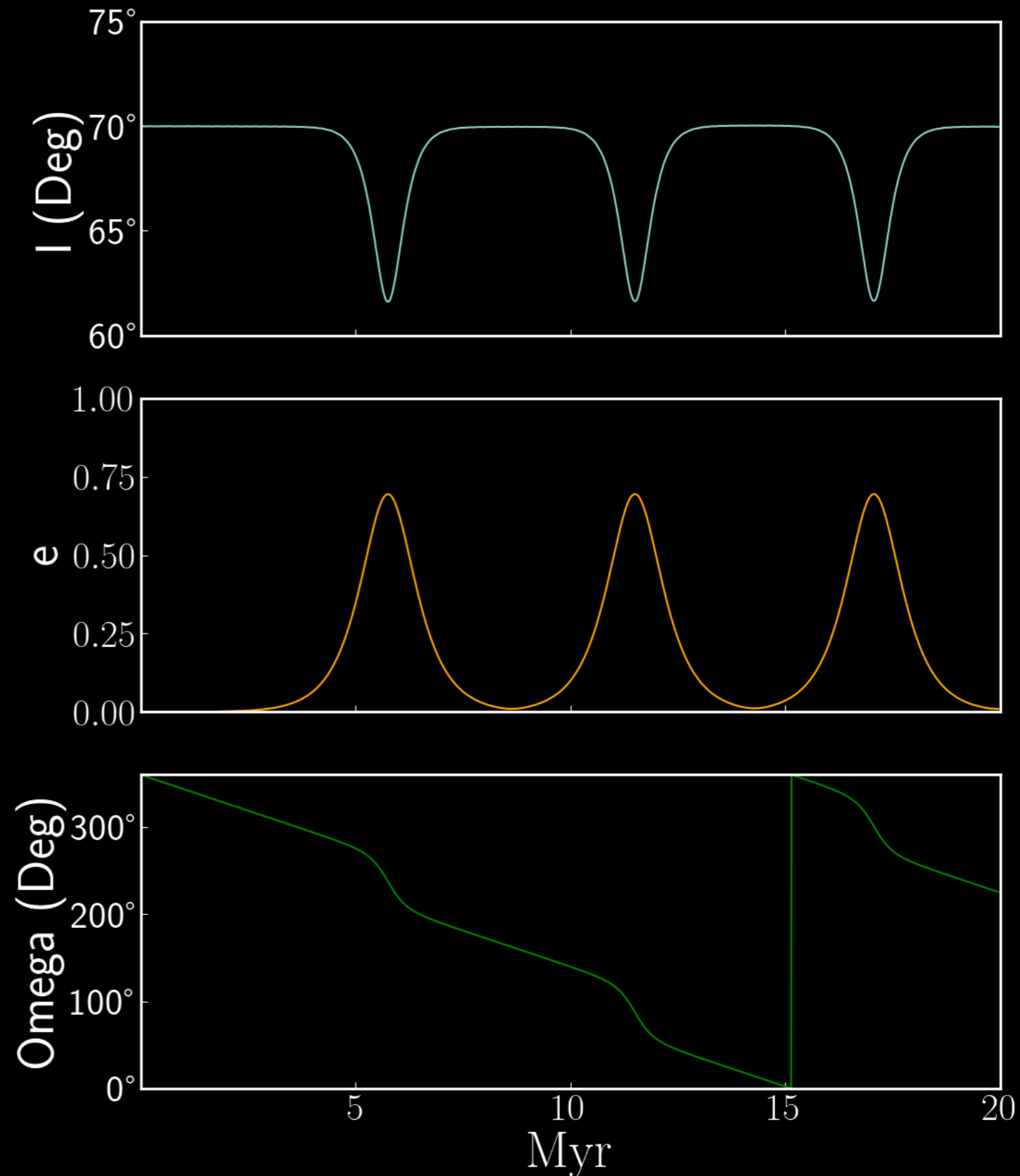
Keplerian Orbital Elements

- Shape of orbit:
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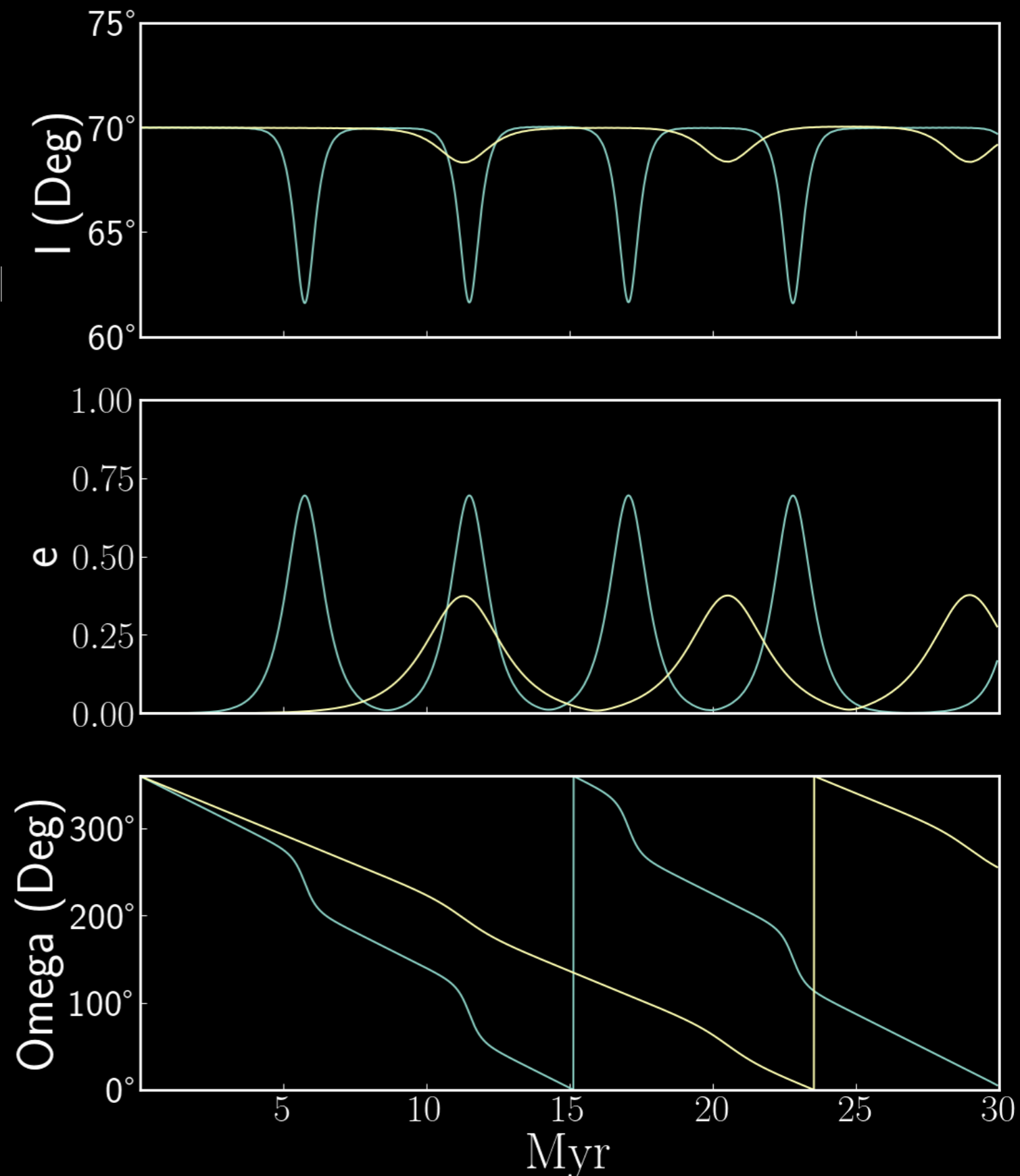
Kozai-Lidov Oscillations

- A binary system perturbed by a massive body.
- Angular momentum of binary is no longer conserved.
- $M_{SMBH} = 4 \cdot 10^6 M_{\odot}$
 $M_{CWD} = 10^4 M_{\odot}$
 $R_{CWD} = 0.1 pc$
 $M_{test} = 10 M_{\odot}$
 $R_{test} = 2.2 \cdot 10^{-2} pc$



Kozai-Lidov Oscillations

- Oscillations damped due to spherically symmetric external potential.
- External potential can be
 - extended stellar cusp.
 - relativistic corrections to newtonian dynamics.
- $M_{SMBH} = 4 \cdot 10^6 M_{\odot}$
 $M_{test} = 10 M_{\odot}$
 $R_{test} = 1.5 \cdot 10^{-2} pc$
 $M_{CWD} = 10^4 M_{\odot}$
 $R_{CWD} = 0.1 pc$



VHS Mechanism

Haas, Šubr & Vokrouhlický (2011)

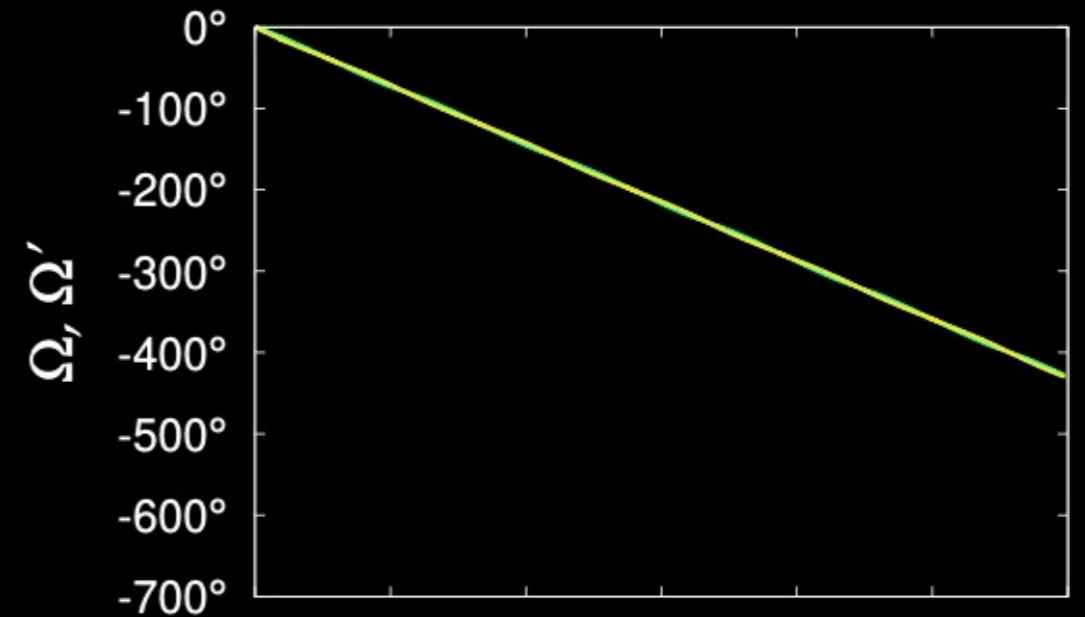
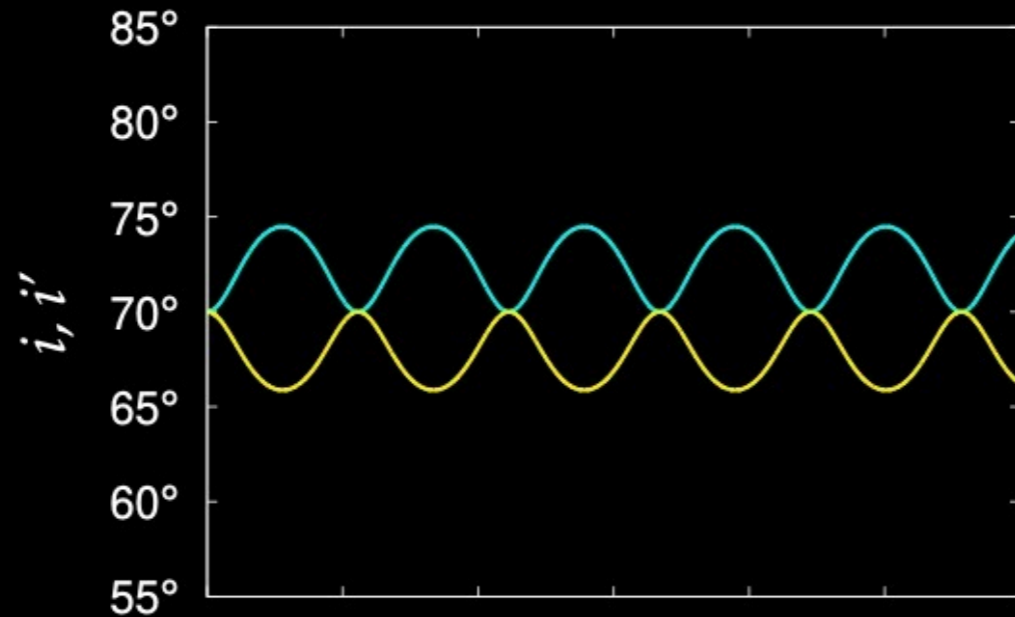
- Four body system
 - Central massive body ($M_{SMBH} = 3.5 \cdot 10^6 M_{\odot}$)
 - 1 massive perturber on circular orbit. ($M_{CND} = 0.3 M_{SMBH}$, $R_{CND} = 1.5 pc$)
 - 2 light bodies on circular orbits:
 - $a_1 = 0.04 R_{CND}$, $a_2 = 0.05 R_{CND}$
 - $e_1 = e_2 = 0$, $i_1 = i_2 = 70^\circ$

	Strong Interaction	Weak Interaction
m_1	$9 \cdot 10^{-6} M_{SMBH}$	$5 \cdot 10^{-6} M_{SMBH}$
m_2	$9 \cdot 10^{-6} M_{SMBH}$	$5 \cdot 10^{-6} M_{SMBH}$

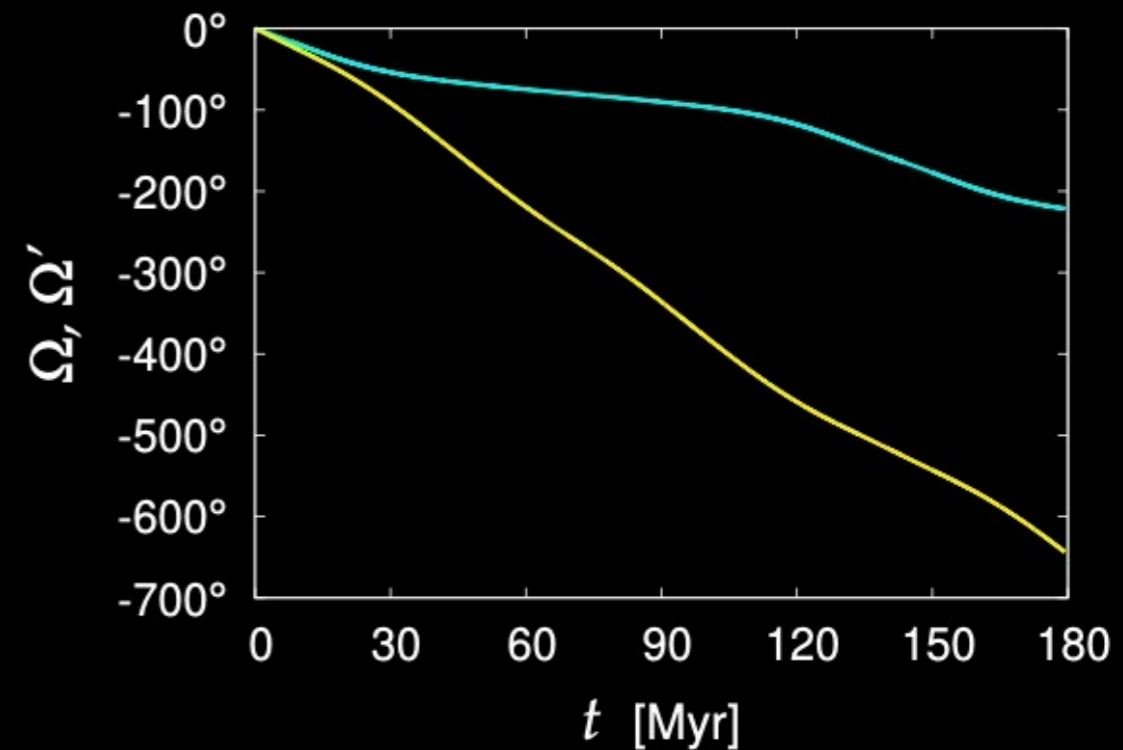
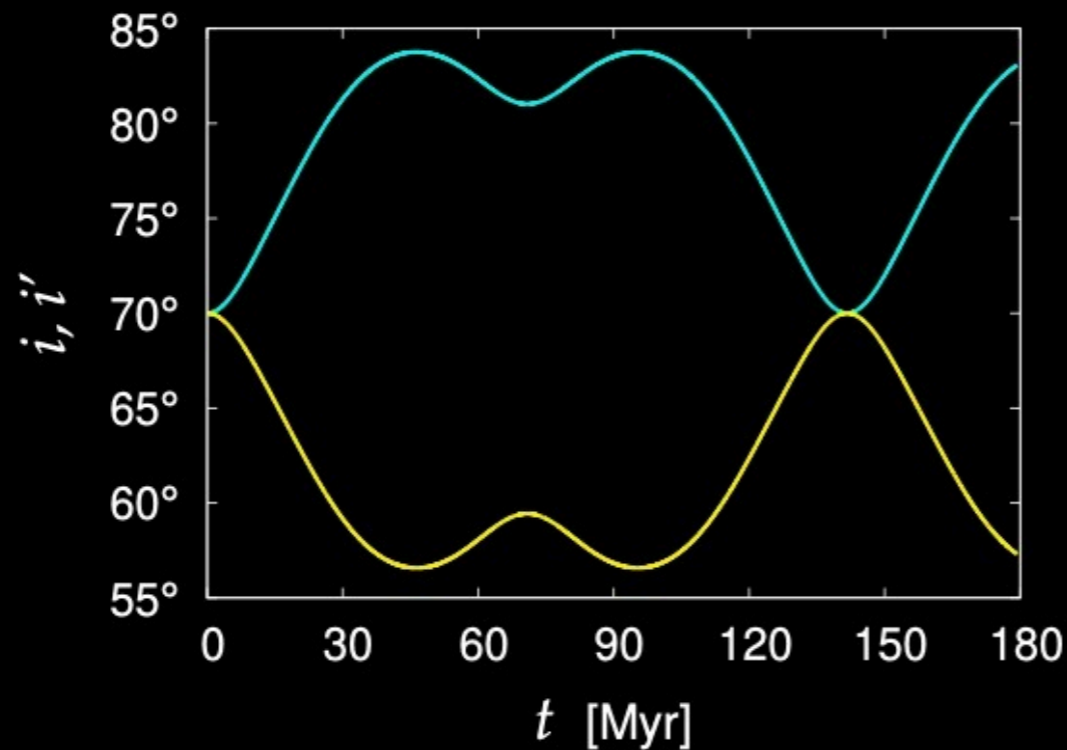
- Spherical external potential from stellar cusp to damp KL oscillations.

VHS Mechanism

Strong
Interaction



Weak
Interaction



VHS Mechanism

External Potential

In Hass et. al. (2011) the averaged external potential to dampen the KL oscillations was given by

$$\overline{\mathcal{R}_c} = -\frac{GmM_c}{\beta R_{CWD}} \left(\frac{a}{R_{CWD}} \right)^\beta \mathcal{I}(e, \beta)$$

where

$$\mathcal{I}(e, \beta) = \frac{1}{\pi} \int_0^\pi (1 - e \cos u)^{1+\beta} du = 1 + \sum_{n \geq 1} a_n e^{2n}$$

And coefficients are given by

$$\frac{a_{n+1}}{a_n} = \left[1 - \frac{3 + \beta}{2(n + 1)} \right] \left[1 - \frac{2 + \beta}{2(n + 1)} \right]$$

And $a_1 = \beta(1 + \beta)/4$ where $\beta = 1/4$

VHS Mechanism with relativistic corrections

External Potential

Rubincam (1977) provides with an external potential

$$V_{GR} = -\frac{GM \cdot h^2}{c^2 r^3}$$

Which can be used to get 1st order corrections. Averaging it like before we get

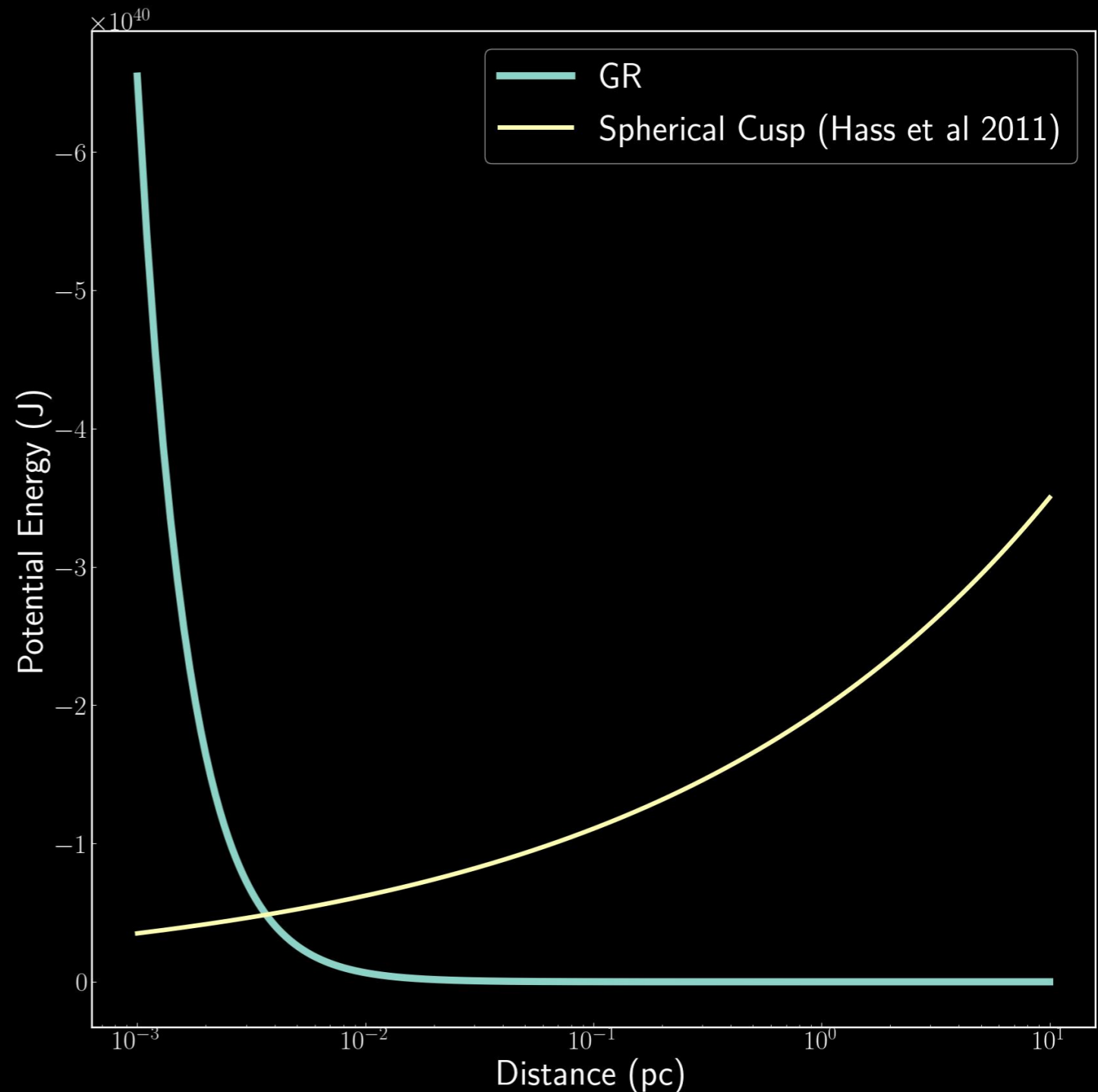
$$\overline{\mathcal{R}}_{GR} = -\frac{GM \cdot mh^2}{c^2 a^3} \mathcal{J}(e, \beta)$$

where $\mathcal{J}(e, \beta)$ is the same but with $\beta = -3$

VHS Mechanism with relativistic corrections

External Potential

- GR correction potential is much stronger in close region around the SMBH.
- Quickly decays so only damps KL oscillations within a small region.



VHS Mechanism with relativistic corrections

Code

- We use *ARWV*, a N-body integration code which calculates PN corrections upto 2.5 orders (Chassonnery et al. 2019).
- It uses the *ARCHAIN* algorithm developed by Mikkola and Merritt (2006, 2008) to calculate velocity dependent forces.

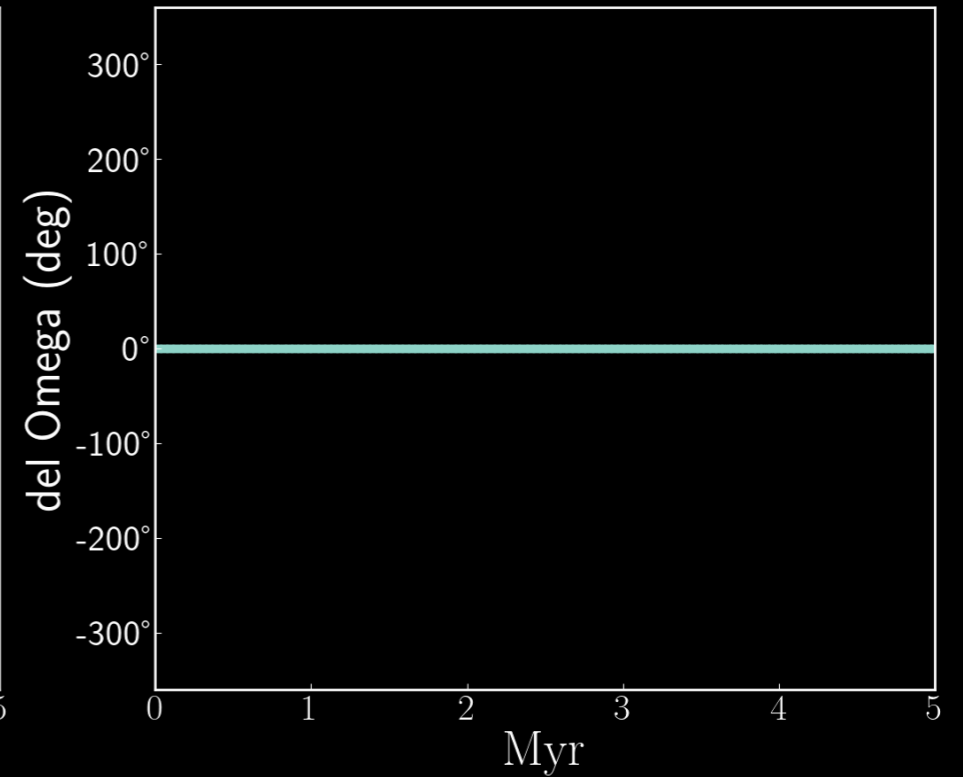
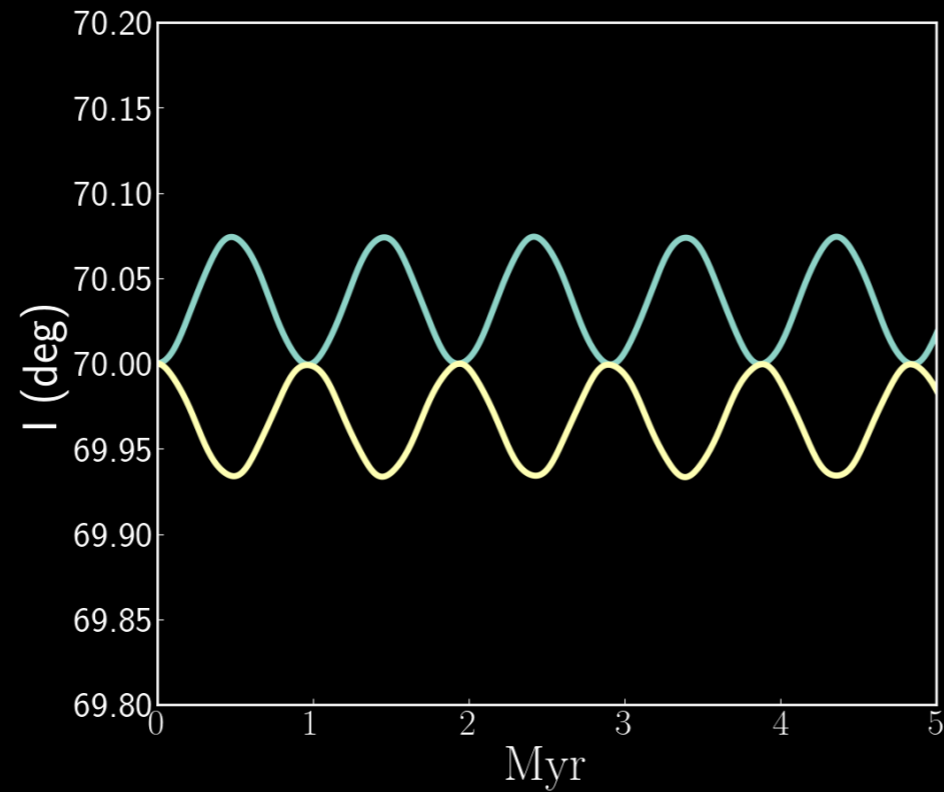
Setup

- Four body system
 - Central massive body ($M_{SMBH} = 4 \cdot 10^6 M_{\odot}$)
 - 1 massive perturber on circular orbit. ($M_{CWD} = 10^4 M_{\odot}, R_{CWD} = 0.1 pc$)
 - 2 light bodies on initially circular orbits:
 - $e_1 = e_2 = 0, i_1 = i_2 = 70^\circ$

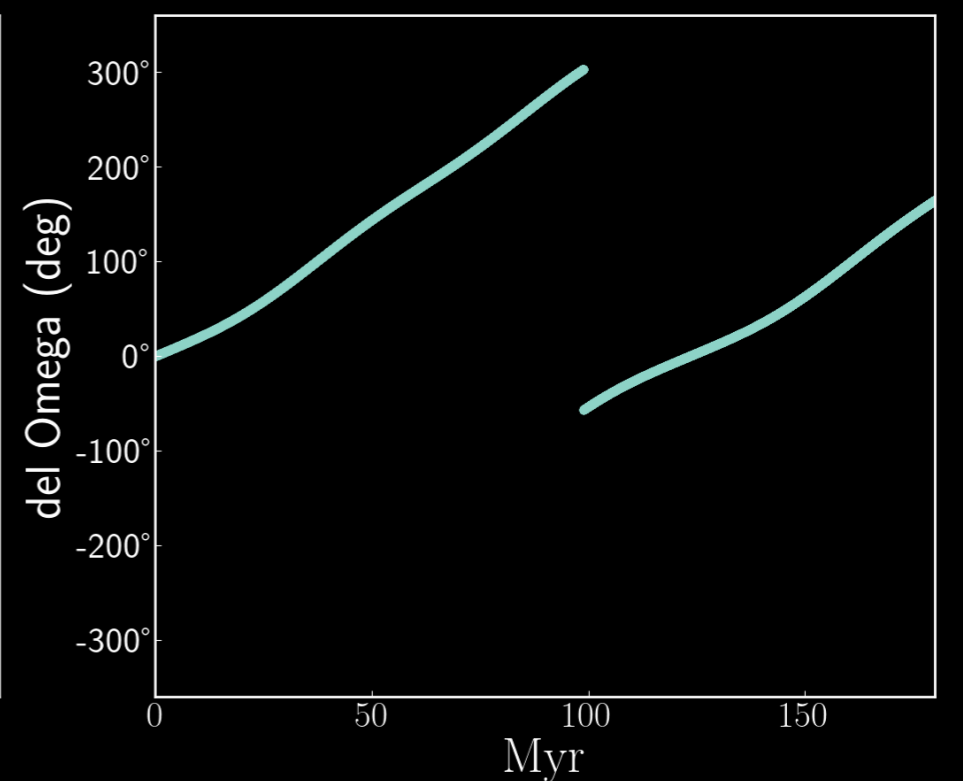
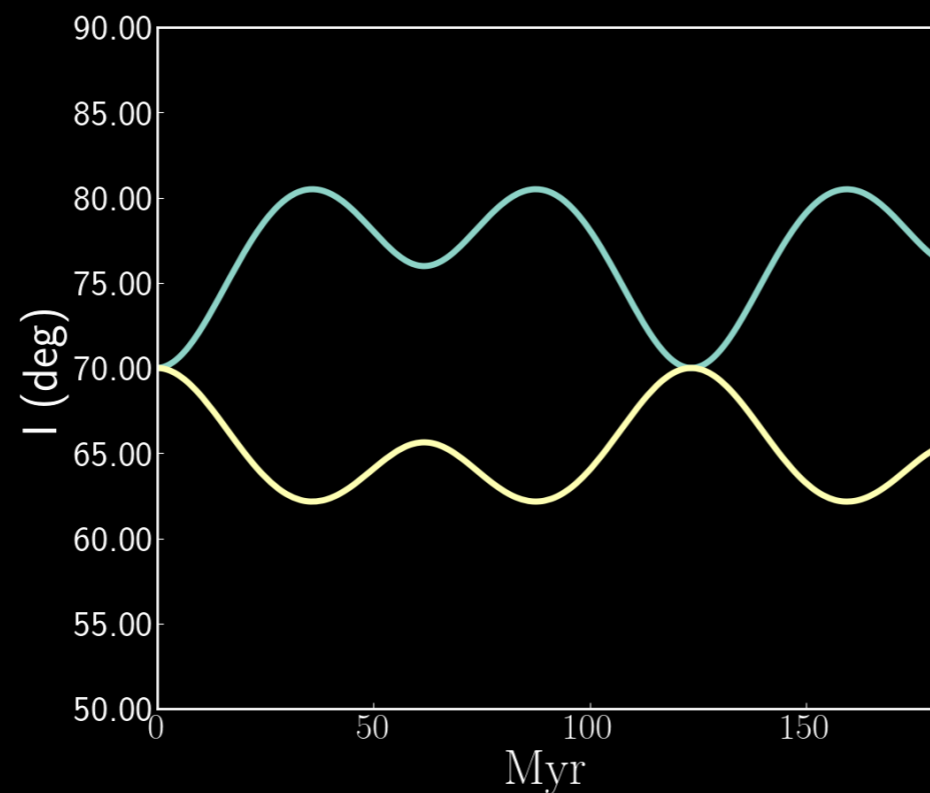
	Strong Interaction	Weak Interaction
m_1	$10M_{\odot}$	$1M_{\odot}$
m_2	$10M_{\odot}$	$1M_{\odot}$
a_1	$3.5 \cdot 10^{-2} R_{CWD}$	$3.5 \cdot 10^{-2} R_{CWD}$
a_2	$4.5 \cdot 10^{-2} R_{CWD}$	$7.0 \cdot 10^{-2} R_{CWD}$

VHS Mechanism with relativistic corrections

Strong
Interaction



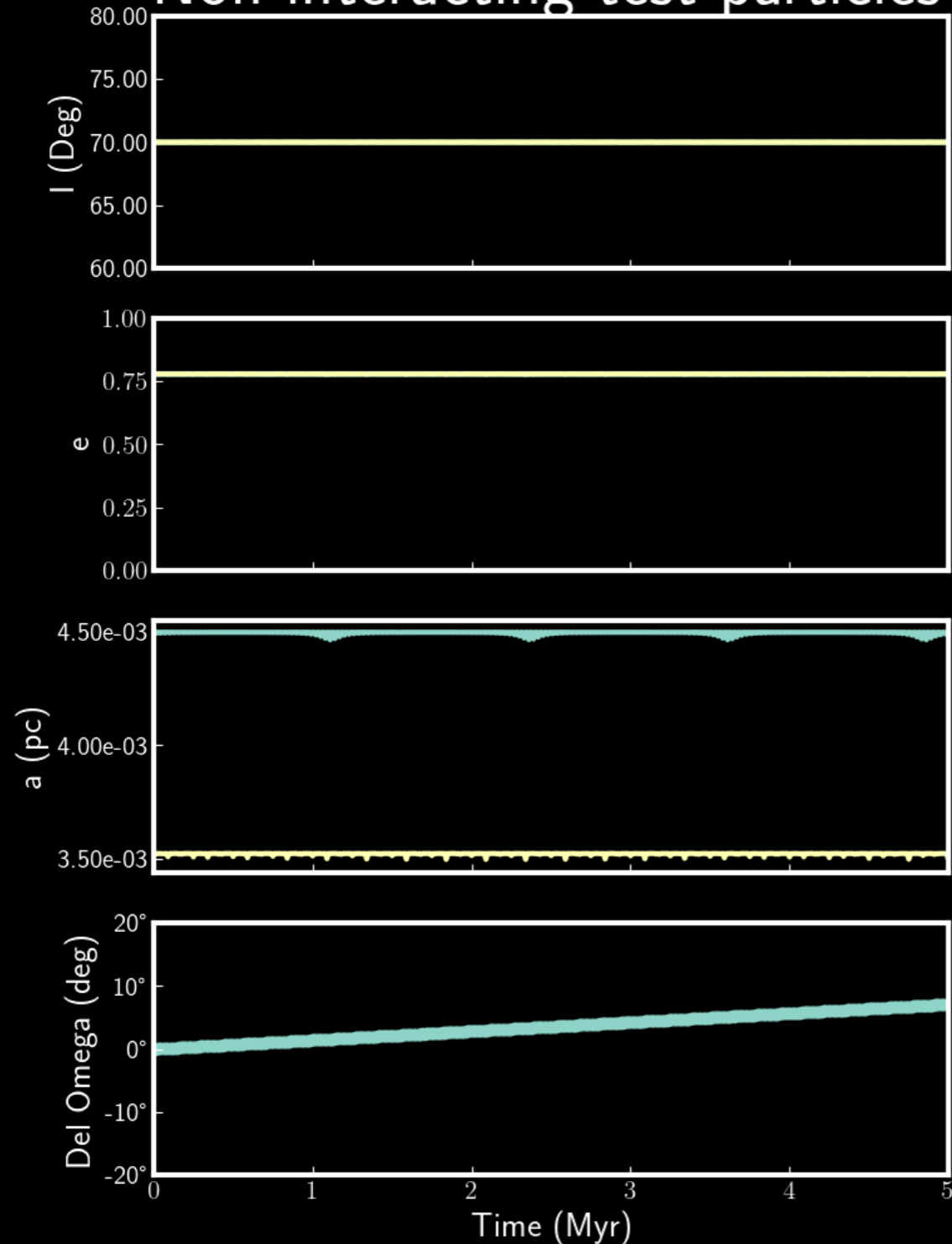
Weak
Interaction



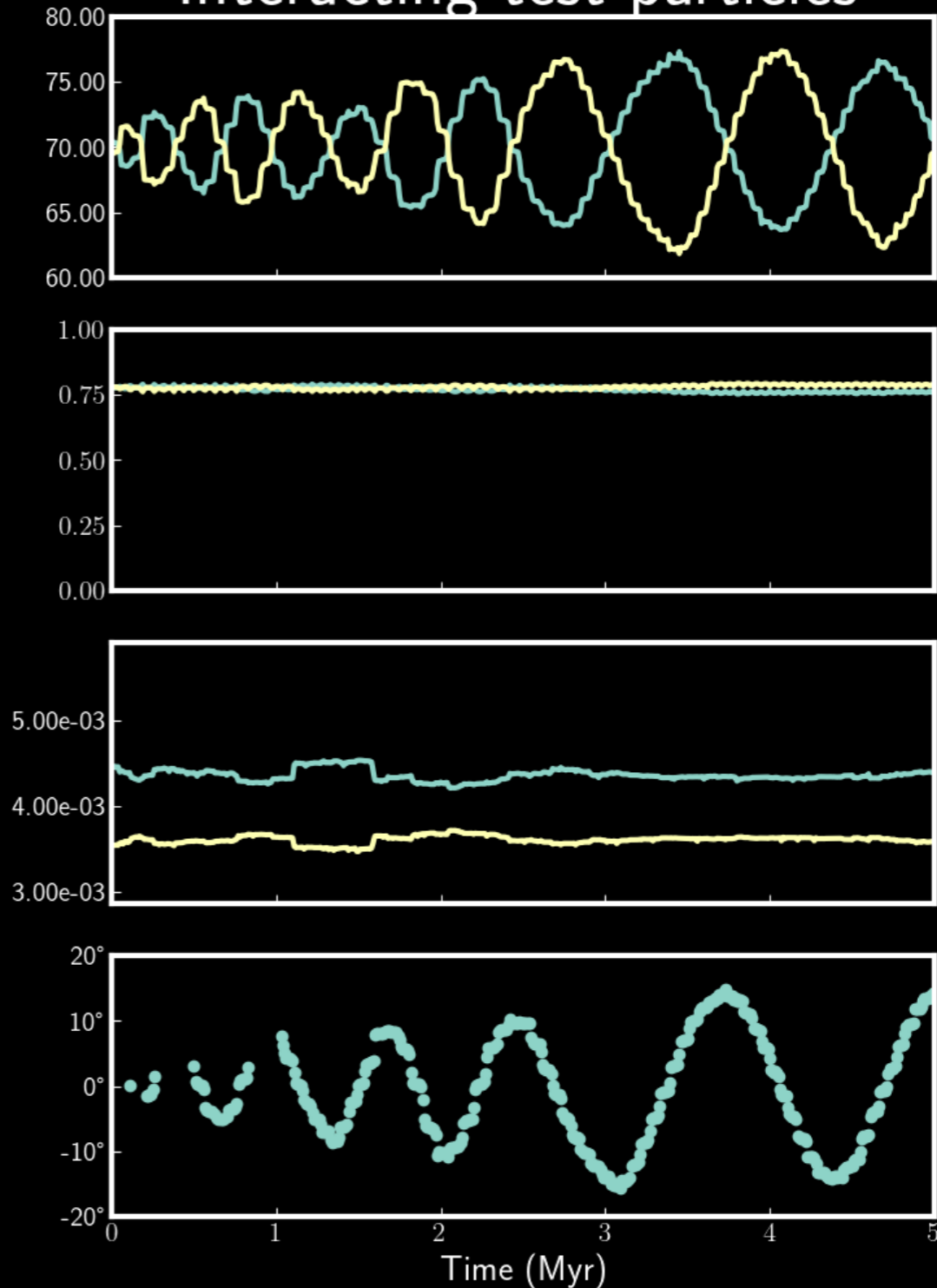
VHS Mechanism with relativistic corrections

Strong Interaction - Non-zero eccentricity

Non-interacting test particles

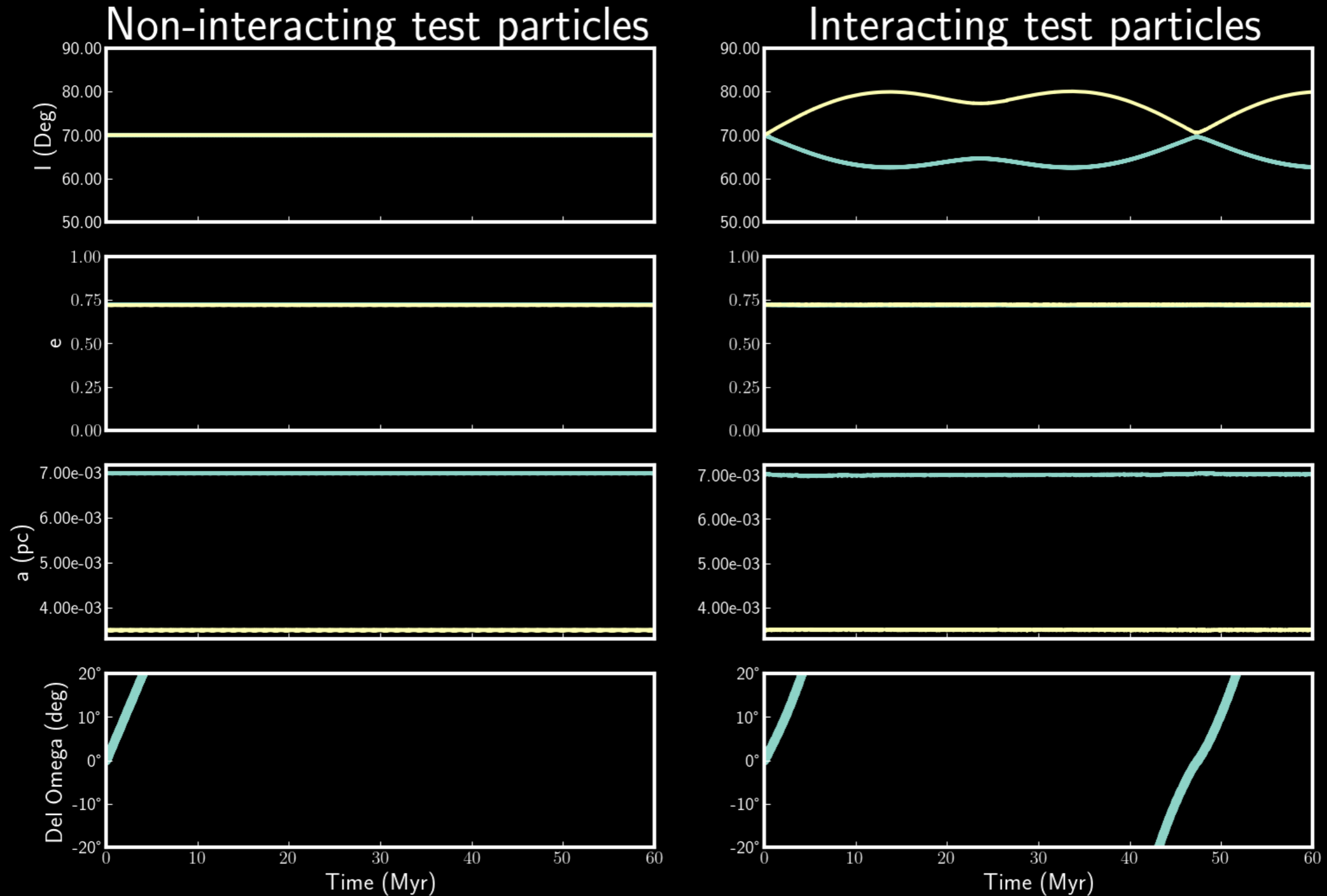


Interacting test particles



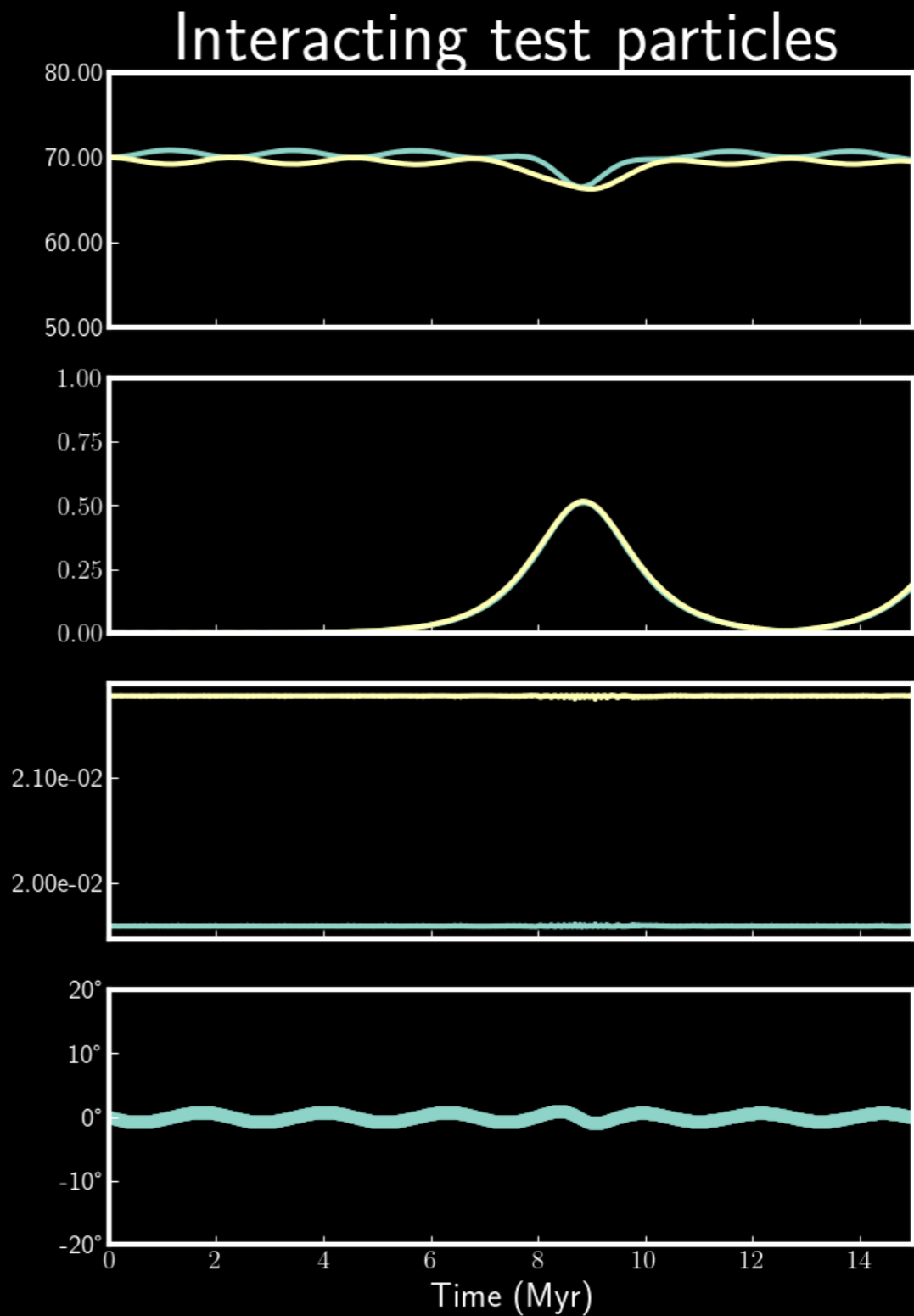
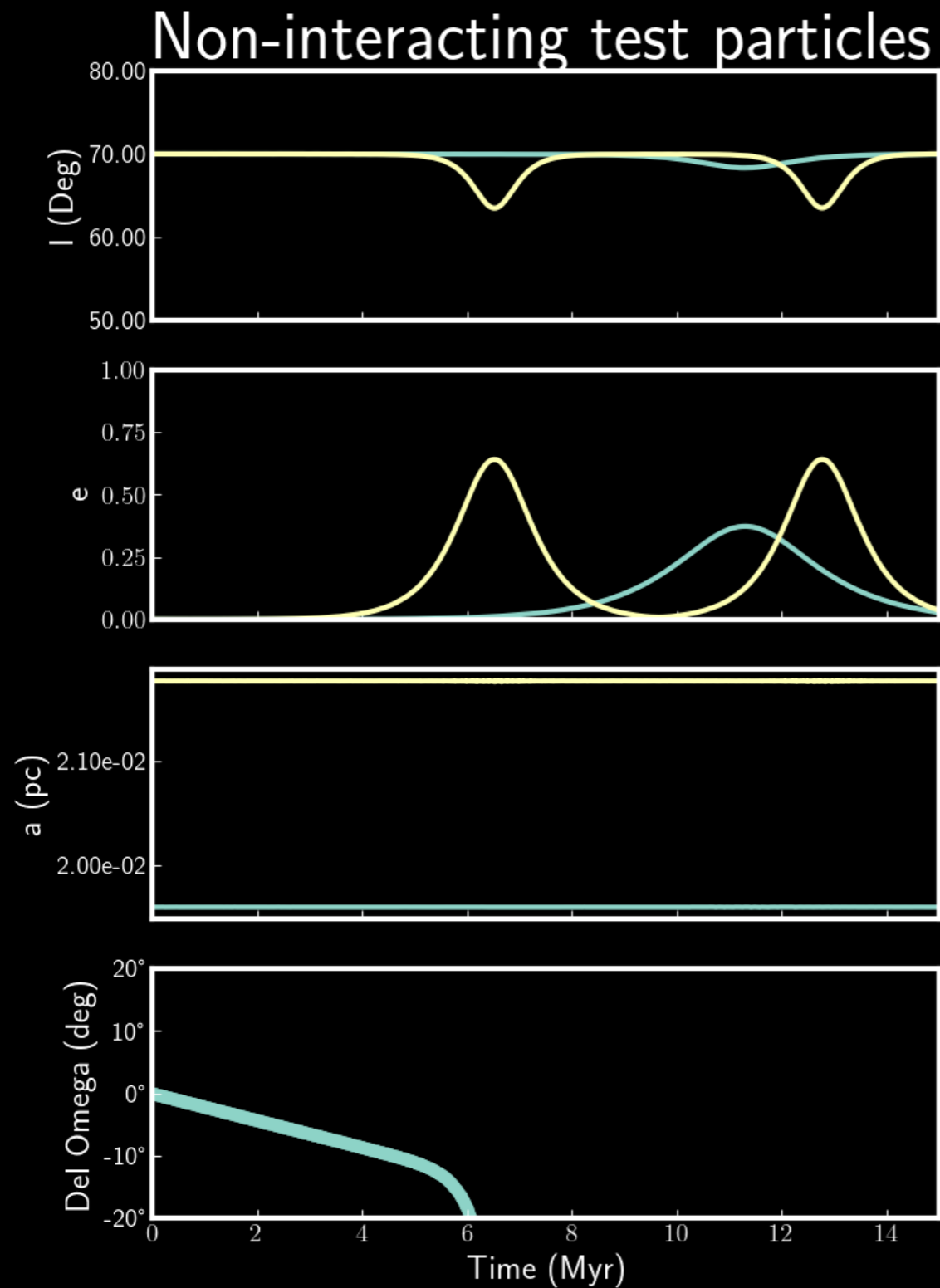
VHS Mechanism with relativistic corrections

Weak Interaction - Non-zero eccentricity



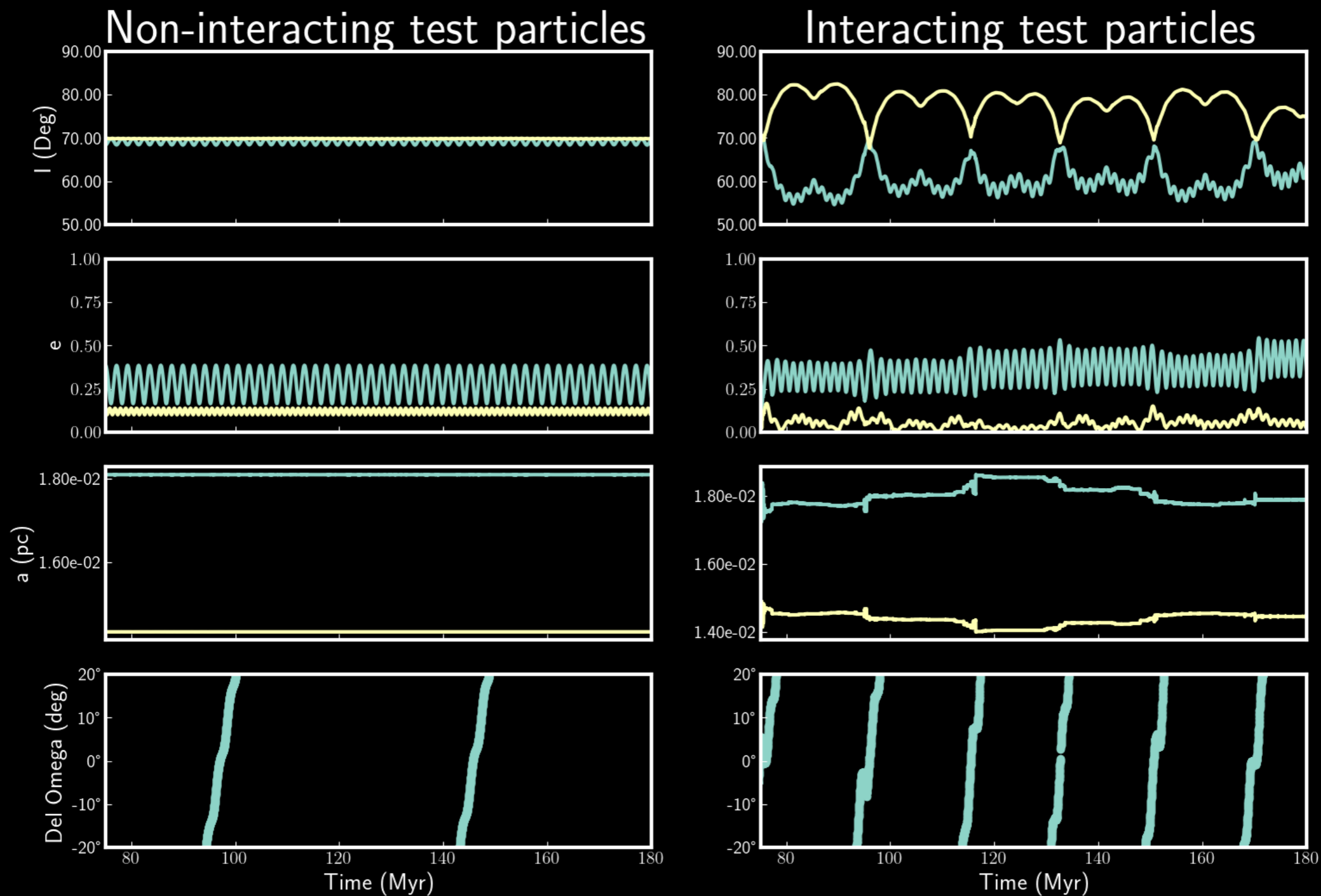
VHS Mechanism with relativistic corrections

Strong Interaction - KL Oscillations



VHS Mechanism with relativistic corrections

Weak Interaction - KL Oscillations



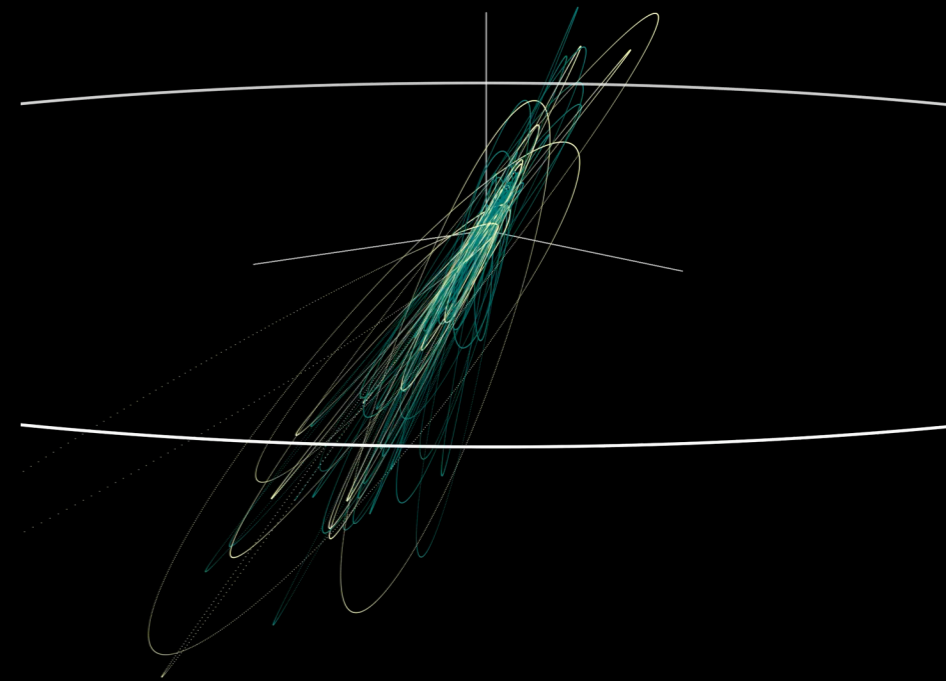
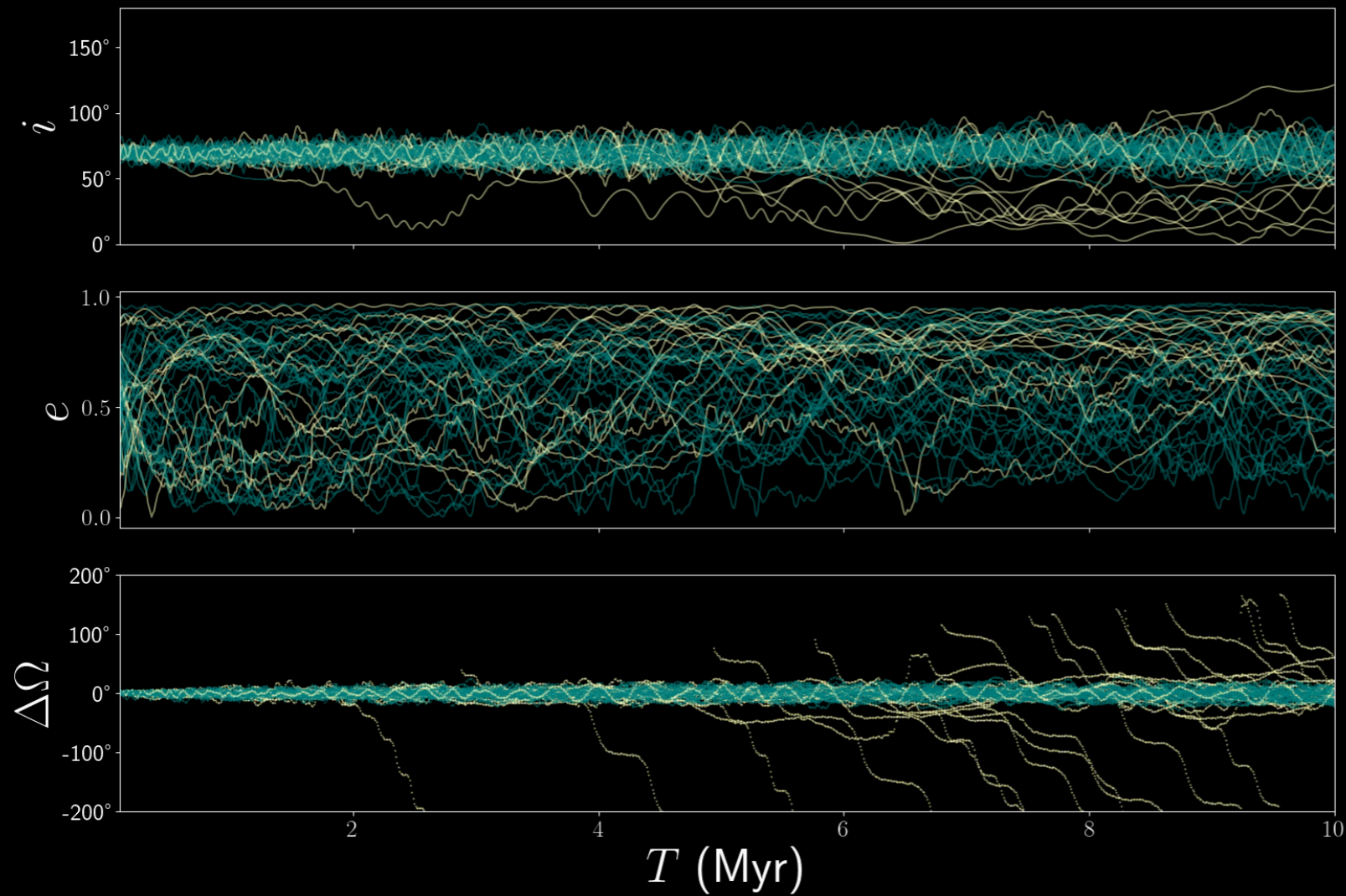
Evolution of a disk of stars

Setup

- Central massive body ($M_{SMBH} = 4 \cdot 10^6 M_{\odot}$)
- 1 massive perturber on circular orbit. ($M_{CWD} = 10^4 M_{\odot}, R_{CWD} = 0.1 pc$)
- Disk of 50 stars:
 - Equal mass of $10M_{\odot}$
 - $e \in (0,1)$
 - $a \in [3.5 \cdot 10^{-4}, 2.0 \cdot 10^{-2}] pc$
 - $i \in [65^{\circ}, 75^{\circ}]$

Evolution of a disk of stars

Evolution



Summary

- The four body dynamics of VHS mechanism are applicable in relativistic regime.
- These dynamics are not only applicable in secular system with damped KL oscillations but
 - can exist in non-eccentric orbits with slight changes.
 - can co-exist with KL oscillations and bind the oscillation together in case of strong interaction.
- These relativistic corrections are applicable to stars in close orbit around Sagittarius A* and these dynamics could be present in that system.